

NONSTATIONARY STAGE OF QUASI-CHERENKOV BEAM INSTABILITY IN PERIODICAL STRUCTURES

K. BATRAKOV¹ and S. SYTOVA²

Institute for Nuclear Problems, Belarusian State University

11 Bobruiskaya, 220050, Minsk, Belarus

E-mail: {¹batrakov, ²sytova}@inp.minsk.by

This contribution is devoted to further investigation by methods of mathematical modeling of quasi-Cherenkov beam instability in periodical structures. Such physical mechanism serves as a basis for Volume Free Electron Lasers (VFEL) that are progressing rapidly. Let us consider the following system of equations describing nonlinear stage of quasi-Cherenkov instability:

$$\frac{\partial E}{\partial t} + a_1 \frac{\partial E}{\partial z} + b_{11}E + b_{12}E_\tau = \Phi \int_0^{2\pi} \frac{2\pi - p}{8\pi^2} (\exp(-i\Theta(t, z, p)) + \exp(-i\Theta(t, z, -p))) dp,$$

$$\frac{\partial E_\tau}{\partial t} + a_2 \frac{\partial E_\tau}{\partial z} + b_{21}E + b_{22}E_\tau = 0,$$

$$\frac{d^2\Theta(t, z, p)}{dz^2} = \Psi \left(k - \frac{d\Theta(t, z, p)}{dz} \right)^3 \operatorname{Re} (E(t - z/u, z) \exp(i\Theta(t, z, p))),$$

$$E|_{z=0} = E_0, \quad E_\tau|_{z=L} = E_1, \quad E|_{t=0} = 0, \quad E_\tau|_{t=0} = 0, \quad \Theta(t, 0, p) = p, \quad \frac{d\Theta(t, 0, p)}{dz} = k - \omega/u,$$

where $t > 0$, $z \in [0, L]$, $p \in [-2\pi, 2\pi]$. This system is a system of integro-differential equations. In addition to temporal argument we have two independent arguments: spatial coordinate z and initial electron phase p . Amplitudes of electromagnetic fields $E(t, z)$, $E_\tau(t, z)$ and coefficients a , b and Φ are complex-valued. Function $\Theta(t, z, p)$ describes phase of electron beam relative to electromagnetic wave. Θ and coefficient Ψ are real. k is a projection of wave vector on axis z . ω is a field frequency. u is an initial electron beam velocity.

Introducing uniform grids on t , z and p

$$\omega_t = \{t_l = lh_t, \quad l = 0, 1, \dots\}, \quad \omega_z = \{z_m = mh_z, \quad m = 0, 1, \dots, M, \quad Mh_z = L\},$$

$$\omega_p = \{p_j = h_p j, \quad j = -N, \dots, -1, 0, 1, \dots, N, \quad h_p N = 2\pi\},$$

we can write the following difference scheme:

$$\widehat{\Theta}_{zz}^j = \Psi \left(k - \widehat{\Theta}_z^j \right)^3 \operatorname{Re} \left(\widetilde{E} \exp(i\Theta^j) \right), \quad \text{where } \widetilde{E} = E(t_l - \alpha t_m, z_m), \quad \alpha = h_z/(h_t u),$$

$$E_t + a_1 \widehat{E}_z + b_{11} \widehat{E} + b_{12} \widehat{E}_\tau = \Phi \sum_{j=0}^N c_j \left(\exp(-i\widehat{\Theta}^j) + \exp(-i\widehat{\Theta}^{-j}) \right), \quad E_{\tau t} + a_2 \widehat{E}_{\tau z} + b_{21} \widehat{E} + b_{22} \widehat{E}_\tau = 0.$$

c_j are coefficients of quadrature formula. We take entier of α . When $t_l - \alpha t_m < 0$, $\widetilde{E} = E(0, z_m)$.

Results of numerical experiments carried out are discussed. In Bragg geometry of VFEL the regime of bifurcations was investigated in detail.